$$
\begin{equation*}
\Delta G^{\circ \prime}=2.3 R T\left[\mathrm{p} K_{\mathrm{a}}^{\mathrm{A}}-\mathrm{p} K_{\mathrm{a}}^{\mathrm{B}}\right] \tag{41}
\end{equation*}
$$

The combination of eq 40 and 41 with the Marcus eq 26 yields

$$
\begin{equation*}
\log \kappa=A+B\left[p K_{\mathrm{a}}^{\mathrm{B}}\right]+C\left[\mathrm{p} K_{\mathrm{a}}^{\mathrm{B}}\right]^{2} \tag{42}
\end{equation*}
$$

where $A=-(R T \ln 10)^{-1}\left\{w^{r}+\Delta G^{\ddagger}{ }_{0}+\left[\mathrm{p} K_{\mathrm{a}}{ }^{\mathrm{A}}\right](R T \ln 10) / 2+\right.$ $\left.\left[\mathrm{p} K_{\mathrm{a}}{ }^{\mathrm{A}}\right]^{2}(R T \ln 10)^{2} / 16 \Delta G^{*} 0\right\rangle+\ln Z / n+7 \Delta E / R T, B=(R T \ln 10)^{-1}-$ $\left[\mathrm{p} K_{\mathrm{a}}^{\mathrm{A}}\right] / 8 \Delta G^{*}{ }_{0}+1 / 2$, and $C=-(R T \ln 10) / 16 \Delta G^{*}{ }_{0}$. The solution of eq 42 was carried out by a polynomial regression analysis on a Digital Equipment Corp. 11/23 computer using a MINC basic subroutine in which $\log \kappa$ was the dependent variable and $\left[\mathrm{p} K_{\mathrm{a}}{ }^{\mathrm{B}}\right.$ ] the independent variable. The second-order coefficient $C=-0.013154$ yielded the values of $\Delta G^{*}$, the first-order coefficient $B=0.551694$ yielded the value of $\left[\mathrm{p} K_{\mathrm{a}}^{\mathrm{A}}\right.$ ], and the zero-order term $A=-7.35117$ yielded the value of $w^{T}$ with a correlation coefficient of 0.995 for all the data in Table $V$ excluding 2,6 -di-tert-butylpyridine (see text).

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Registry No. $\mathrm{Fe}(\text { phen })_{3}\left(\mathrm{PF}_{6}\right)_{3}, 28277-57-8 ; \mathrm{Dz}, 7782-39-0$; hexamethylbenzene, 87-85-4; pentamethylbenzene, 700-12-9; 1,2,4,5-tetramethylbenzene, 95-93-2; 1,2,3,4-tetramethylbenzene, 488-23-3; 1,2,3,5tetramethylbenzene, 527-53-7; hexamethylbenzene cation radical, 34473-51-3; 1,2,4,5-tetramethylbenzene cation radical, 34473-49-9; pentamethylbenzene cation radical, 34473-50-2; 1,2,3,4-tetramethylbenzene cation radical, 34528-28-4; 2-fluoropyridine, 372-48-5; 2chloropyridine, 109-09-1; 3-pyridinecarbonitrile, 100-54-9; 4-pyridinecarbonitrile, 100-48-1; 3-chloropyridine, 626-60-8; 3-fluoropyridine, 372-47-4; pyridine, 110-86-1.

# Reaction Mechanisms of Oxidative Addition [ $\mathrm{H}_{2}+$ $\mathrm{Pt}^{0}\left(\mathrm{PH}_{3}\right)_{2} \rightarrow \mathrm{Pt}^{\mathrm{II}}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ ] and Reductive Elimination $\left[\mathrm{Pt}^{\mathrm{II}}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2} \rightarrow \mathrm{CH}_{4}+\mathrm{Pt}^{0}\left(\mathrm{PH}_{3}\right)_{2}\right]$. Ab Initio MO Study 

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#### Abstract

Reaction mechanisms of the oxidative addition of $\mathrm{H}_{2}$ to two-coordinate $\mathrm{Pt}^{0}\left(\mathrm{PH}_{3}\right)_{2}$ and the reductive elimination of $\mathrm{CH}_{4}$ from four-coordinate $\mathrm{Pt}^{11}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ are studied by ab initio RHF and CI calculations with the energy gradient method within the framework of the relativistic effective core potential approximation for Pt core electrons. Fully optimized geometries of transition states of both reactions as well as the reactants and products have been obtained. The fact that the oxidative addition more commonly takes place for $\mathrm{H}_{2}$ and only the reductive elimination does for $\mathrm{CH}_{4}$ can be explained in terms of calculated exothermicity. The $\mathrm{H}_{2}$ oxidative-addition reaction is suggested to pass through an early transition state that would lead directly to a cis product and then to be pushed toward a trans product by a steric repulsion between bulky phosphine ligands. A large deuterium kinetic isotope effect experimentally found in the reductive elimination reaction of $\mathrm{CH}_{3} \mathrm{D}$ from $\mathrm{Pt}(\mathrm{D})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PPh}_{3}\right)_{2}$ is accounted for in terms of the calculated four-coordinate transition state, where the reaction coordinate is CPtH bending. A decrease in the interligand angle has been found to increase the reactivity of the metal center by selectively activating one of the d orbitals. Both donation and back-donation between the metal and $\mathrm{H}_{2}$ have been found to be important at the transition state of the $\mathrm{H}_{2}$ addition reaction.


The oxidative addition of molecular hydrogen and aliphatic RH and the corresponding reverse reductive elimination are important reactions as elementary processes in many catalytic and synthetic reactions. ${ }^{2}$ Recent studies by Otsuka, Yoshida, and their colleagues on preparation and reaction of $\mathrm{H}_{2}$ with two-coordinate platinum ( 0 ) [and palladium(0)] phosphine complexes present an interesting chemistry of homogeneous catalytic activities. ${ }^{3-5}$ Some of them easily absorb molecular hydrogen ${ }^{4}$ at room temperature to give dihydride complexes, which are identified by NMR and IR studies to be trans products. Oxidative-addition products have been found to be diamagnetic, indicating a closed-shell singlet ground state, which is reasonable for $\mathrm{d}^{8}$ square-planar complexes. ${ }^{6}$ Reversible reactions with $\mathrm{H}_{2}$ at room temperature ${ }^{5}$ are also found for some of $\mathrm{Pt}(0)$-chelating phosphine complexes. The symmetry

[^0]rule suggests that the oxidative addition of $\mathrm{H}_{2}$ occurs in a cis fashion; ${ }^{2}$ however, a trans mechanism has been proposed for the addition of $\mathrm{H}_{2}$ to Ir complexes. ${ }^{?}$
The reductive-elimination reaction of $\mathrm{CH}_{4}$ from $\mathrm{Pt}(\mathrm{H})$ $\left(\mathrm{CH}_{3}\right)\left(\mathrm{PPh}_{3}\right)_{2}$ has been found by Abis, Sen, and Halpern. ${ }^{8}$ The reaction rate was unaffected by the presence of excess phosphine ligands, suggesting strongly that the elimination proceeds through a four-coordinate transition state. They observed a large deuterium kinetic isotope ratio of 3.3 for the decomposition reaction of $\mathrm{Pt}(\mathrm{D})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PPh}_{3}\right)_{2}$. On the other hand, small isotope ratios of 1.2-1.3 have been observed for $\mathrm{H}_{2}$ and RH oxidative-addition reactions. ${ }^{9,10}$ The oxidative addition of aliphatic CH is considered to occur in a $\mathrm{H}-\mathrm{D}$ exchange reaction of alkane with $\mathrm{K}_{2} \mathrm{PtCl}_{4}{ }^{11}$ as a catalyst and in an internal metalation ${ }^{12}$ where a CH bond in a ligand such as $\mathrm{PR}_{3}$ is broken, although in both of the reactions the RH adduct has not been isolated. However, aromatic and olefinic CH's are known to add oxidatively to $\mathrm{Ir}, \mathrm{Pd}, \mathrm{Fe}, \mathrm{Ru}$, and Os compounds. ${ }^{13}$

[^1]Many reductive-elimination reactions of $\mathrm{RR}^{\prime}$ have been found for $\mathrm{Ni}, \mathrm{Pd}$, and Pt compounds. ${ }^{14}$ The oxidative addition of $\mathrm{RR}^{\prime}$, however, has not been found, except for metallacycle formation with strained ring hydrocarbons such as a reaction of cubane with a rhodium compound ${ }^{15}$ and a reaction of cyclopropane with a platinum compound. ${ }^{16}$

From a study of $\operatorname{Pd}\left[\operatorname{PPh}(t-\mathrm{Bu})_{2}\right]_{2}$ and $\operatorname{Pt}\left[\operatorname{PPh}(t-\mathrm{Bu})_{2}\right]_{2}$, where $\mathrm{O}_{2}$ coordination is reversible for the former while it is irreversible for the latter, Otsuka et al. have suggested the difference in an interligand angle as an origin of this reactivity difference and proposed to control the reactivity of metal phosphine complexes with the interligand angle as well as the steric size and basicity of phosphine ligands. ${ }^{17,18}$

A few theoretical studies have been published on these reactions. Hoffmann et al. studied the reductive elimination of alkane from dialkyl $\mathrm{Ni}, \mathrm{Pd}, \mathrm{Pt}$, and Au compounds by the extended Hückel method. ${ }^{19}$ The oxidative addition of $\mathrm{H}_{2}$ to $\mathrm{RhCl}\left(\mathrm{PH}_{3}\right)_{3}$ has been studied by Dedieu and Strich with the extended Hücke ${ }^{20}$ and the ab initio method. ${ }^{21}$ They examined possible schemes of hydrogenation of olefins catalyzed by rhodium compounds. ${ }^{21}$ The bonding interaction between the $\mathrm{H}_{2}$ antibonding orbital and an occupied Rh d orbital has been suggested to facilitate HH bond breaking in this reaction as well as in the $\left(\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2} \mathrm{Mo}+\mathrm{H}_{2}$ reaction. ${ }^{22}$ Whitesides et al. ${ }^{23}$ discussed a correlation between the relative rates of reductive elimination ( $\mathrm{L}_{2} \mathrm{MXY} \rightarrow \mathrm{L}_{2} \mathrm{M}+$ $\mathrm{XY} ; \mathrm{L}=\mathrm{PH}_{3} ; \mathrm{M}=\mathrm{Pt}, \mathrm{Pd}$, and $\mathrm{Ni} ; \mathrm{XY}=(\mathrm{H})_{2},(\mathrm{H})\left(\mathrm{CH}_{3}\right)$, and $\left.\left(\mathrm{CH}_{3}\right)_{2}\right)$ and the molecular orbital character of the starting compounds using self-consistent field-X $\alpha$-scattered-wave calculations.

Noell and Hay have recently studied the oxidative-addition reactions $\mathrm{H}_{2}+\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2} \rightarrow$ cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{H}_{2}+\mathrm{Pt}([\mathrm{P}-$ $\left.\left(\mathrm{CH}_{3}\right)_{3}\right]_{2} \rightarrow$ cis- $\mathrm{Pt}(\mathrm{H})_{2}\left[\mathrm{P}\left(\mathrm{CH}_{3}\right)_{3}\right]_{2}$ by ab initio HF, MCSCF, and CI methods with a relativistic effective core potential and a flexible Gaussian basis set for the valence electrons. ${ }^{24}$ They optimized three important geometrical parameters within $C_{2 v}$ symmetry by the energy method and determined relative energies of reactants, products, and transition states. They obtained an exothermicity of $7 \mathrm{kcal} / \mathrm{mol}$ and an activation barrier of $17 \mathrm{kcal} / \mathrm{mol}$ and
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Scheme I

therefore a barrier of $24 \mathrm{kcal} / \mathrm{mol}$ for the reverse reductive elimination.

Recently the energy gradient method in the ab initio MO method has become a most powerful tool for exploring potential energy surfaces. ${ }^{25}$ In the method the first derivative, i.e., the gradient of the electronic energy with respect to nuclear coordinates is evaluated analytically and efficiently. The gradient is then used to locate the equilibrium and the transition-state geometry in which all the geometrical parameters are fully optimized. The force constant matrix is evaluated by numerical differentiation of the gradient. The energy gradient method has been applied in the $a b$ initio MO method to many reaction systems with several first-row atoms and their transition states have been determined. ${ }^{25}$ However until recently ${ }^{26-28}$ the gradient method has not been used for transition-metal reactions. We have published preliminary results of fully optimized molecular structures of two-coordinate $\mathrm{Pt}(0)$ diphosphine, $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, four-coordinate cis- and trans- $\mathrm{Pt}-$ $(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2},{ }^{26}$ and a transition state of oxidative addition of $\mathrm{H}_{2}$ to $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2} .{ }^{27}$

In this paper we report a full account of our ab initio MO study of the oxidative addition reaction of $\mathrm{H}_{2}$ to $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ as well as the reductive elimination of $\mathrm{CH}_{4}$ from $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$. The purposes of the study are as follows: (1) to check how well the present method reproduces equilibrium geometries and IR normal mode frequencies by comparing them with the experimental results, (2) to elucidate features of potential surfaces for reactions of $\mathrm{H}_{2}$ and $\mathrm{CH}_{4}$ with $\mathrm{Pt}(0)$ diphosphines by determining transi-tion-state geometries and energies and to compare differences between $\mathrm{H}_{2}$ reactions and $\mathrm{CH}_{4}$ reactions, (3) to calculate theoretical values of the deuterium kinetic isotope ratio for the reductive elimination of $\mathrm{Pt}(\mathrm{D})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ and the oxidative addition of $\mathrm{D}_{2}$ to $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, to be compared with experiments, (4) to investigate, using a theoretical model, the effect of the interligand PPtP angle on the reactivity, and (5) to discuss detailed mechanisms of oxidative-addition and reductive-elimination reactions. The second section describes methods of calculation. Results and discussions are given in the third and fourth sections, followed by a summary in the last section.

## Methods of Calculation

We have optimized all the geometrical parameters of $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, cis and trans forms of $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, and the transition states (Scheme I) by using the energy gradient technique ${ }^{26}$ at the restricted Hartree-Fock (RHF) level under the effective core potential (ECP) approximation. ${ }^{29}$ Calculations were performed for a closed-shell singlet state, the ground state. A smaller basis set was used in geometry optimization and normal mode calculations. This basis set consists of valence double (3s 3 p 5 d ) / $[2 \mathrm{~s} 2 \mathrm{p} 2 \mathrm{~d}] \mathrm{set}^{29}$ for $\mathrm{Pt}, 3-21 \mathrm{G}^{30}$ for the hydride H and the methyl group, and STO- $2 \mathrm{G}^{31}$ for $\mathrm{PH}_{3}$. The relativistic ECP obtained by Basch and Topiol ${ }^{29}$ was used to replace core electrons for Pt . The smaller basis set reproduces well the geometries of

[^2]Table I. Mulliken Population Analysis of $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}, \mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, and $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}{ }^{a}$

|  | $\begin{gathered} \text { linear } \\ \operatorname{Pt}\left(\mathrm{PH}_{3}\right)_{2} \end{gathered}$ | trans$\mathrm{Pt}(\mathrm{H})_{2}{ }^{-}$ $\left(\mathrm{PH}_{3}\right)_{2}$ | $\begin{gathered} \text { trans- } \\ \mathrm{Pt}(\mathrm{H})- \\ \left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2} \end{gathered}$ | $\begin{gathered} \text { bent } \\ \mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}{ }^{b} \\ \hline \end{gathered}$ | $\begin{gathered} \text { cis- } \\ \mathrm{Pt}(\mathrm{H})_{2}- \\ \left(\mathrm{PH}_{3}\right)_{2} \end{gathered}$ | $\begin{gathered} c i s- \\ \mathrm{Pt}(\mathrm{H})- \\ \left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2} \end{gathered}$ | transition state |  | $\mathrm{CH}_{4}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  | $\begin{aligned} & \hline \mathrm{Pt}(\mathrm{H})_{2^{-}} \\ & \left(\mathrm{PH}_{3}\right)_{2}{ }^{c} \\ & \hline \end{aligned}$ | $\begin{gathered} \mathrm{Pt}(\mathrm{H})- \\ \left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}{ }^{d} \end{gathered}$ |  |
| Gross ${ }^{\text {e }}$ |  |  |  |  |  |  |  |  |  |
| Pt | 10.52 | 10.67 | 10.39 | 10.23 | 10.67 | 10.43 | 10.40 | 10.39 |  |
| s | 0.96 | 0.81 | 0.73 | 0.43 | 0.91 | 0.83 | 0.62 | 0.59 |  |
| p | 0.18 | 0.89 | 0.74 | 0.36 | 0.76 | 0.67 | 0.41 | 0.56 |  |
| $x$ | 0.01 | 0.01 | 0.03 | 0.02 | 0.02 | 0.04 | 0.01 | 0.02 |  |
| $y$ | 0.02 | 0.67 | 0.48 | 0.07 | 0.39 | 0.27 | 0.17 | 0.26 |  |
| $z$ | 0.15 | 0.21 | 0.23 | 0.28 | 0.35 | 0.36 | 0.23 | 0.28 |  |
| d | 9.38 | 8.98 | 8.92 | 9.43 | 9.00 | 8.93 | 9.37 | 9.24 |  |
| $x y$ | 2.00 | 2.00 | 1.99 | 1.94 | 1.96 | 1.93 | 1.99 | 1.91 |  |
| $x z$ | 1.88 | 1.92 | 1.92 | 1.90 | 1.93 | 1.94 | 1.89 | 1.96 |  |
| $y z$ | 1.88 | 1.93 | 1.93 | 1.87 | 1.33 | 1.90 | 1.82 | 1.73 |  |
| $x^{2}$ | 1.43 | 1.42 | 1.43 | 1.49 | 1.42 | 1.43 | 1.46 | 1.46 |  |
| $y^{2}$ | 1.43 | 0.92 | 0.87 | 1.15 | 1.20 | 0.83 | 1.30 | 1.06 |  |
| $z^{2}$ | 0.77 | 0.78 | 0.78 | 1.08 | 1.17 | 0.90 | 0.90 | 1.11 |  |
| H |  | 1.03 | 1.02 |  | 0.95 | 0.97 | 0.97 | 0.76 | 0.80 |
| $\mathrm{PH}_{3}$ | 7.74 | 7.63 | 7.64 | 7.88 | 7.71 | 7.68 | 7.83 | 7.82, 7.875 |  |
| $\mathrm{CH}_{3}$ |  |  | 9.31 |  |  | 9.24 |  | 9.16 | 9.20 |
| Overlap |  |  |  |  |  |  |  |  |  |
| $\mathrm{Pt}-\mathrm{H}$ |  | 0.38 | 0.38 |  | 0.39 | 0.39 |  | 0.18 |  |
| $\mathrm{Pt}-\mathrm{P}$ | 0.35 | 0.33 | 0.33 | 0.31 | 0.26 | 0.25 |  | $0.34,0.28^{\prime}$ |  |
| $\mathrm{Pt}-\mathrm{C}$ |  |  | 0.29 |  | 0.32 |  |  | 0.13 |  |

${ }^{a}$ The analysis is based on RHF calculation with the larger basis set. The direction of the coordinate axes is shown in Scheme I. ${ }^{b}$ The molecular geometry of the bent $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ is set to be the same as that of the corresponding cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$. ${ }^{c}$ The Pt and two P atoms are on the $y z$ plane, and the symmetry axis is on the $y$ axis. ${ }^{d}$ The Pt and two P atoms are on the $y z$ plane, and the PtH bond is on the $z$ axis. ${ }^{e}$ Core electrons of Pt and $\mathbf{P}$ atoms are not taken into account. Therefore the total gross populations of a free Pt atom and $\mathrm{PH}_{3}$ molecule become 10 and 8 , respectively, ${ }^{f}$ The former value corresponds to the $\mathrm{PH}_{3}$ adjacent to the hydride.
$\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, and trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ obtained by the larger basis set. ${ }^{26}$ The larger set is obtained from the smaller set by replacing basis functions for H by 21 G and P by LP-31G with the nonrelativistic ECP approximation. ${ }^{32}$ In calculations of relative energies of reactants, products, and transition states, configuration interaction calculations including all the single and double excitations of valence electrons relative to the RHF configuration, ${ }^{33}$ often called SD-CI, were carried out with the larger basis set at the geometries optimized at the smaller basis RHF level. The number of configurations included a mounts to be about 67000 and 150000 for $\mathrm{Pt}\left(\mathrm{H}_{2}\right)_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{Pt}(\mathrm{H})(\mathrm{C}-$ $\left.\mathrm{H}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, respectively, in the $C_{s}$ symmetry. The correction for unlinked quadruple excitations, called QC, has also been included on the basis of Davidson's formula. ${ }^{34}$

We used IMSPAK ${ }^{35}$ for ECP geometry optimization and the direct CI program in alchemy for CI calculation. For d functions $x^{2}+y^{2}+z^{2}$ components are excluded from RHF and CI calculations. Errors caused by the ECP approximation had been examined for geometries and relative energies of the first- and second-row polyatomic molecules. ${ }^{36}$
$\mathbf{P t}\left(\mathrm{PH}_{3}\right)_{2}$, cis $-\mathbf{P t}(\mathrm{H})(\mathbf{R})\left(\mathrm{PH}_{3}\right)_{2}$, and trans $-\mathbf{P t}(\mathrm{H})(\mathbf{R})\left(\mathrm{PH}_{3}\right)_{2}$ ( $\mathrm{R}=\mathrm{H}$ or $\mathrm{CH}_{3}$ ). (A) Electronic Structures. The calculated orbital energy levels of $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, trans- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, and cis- $\mathrm{Pt}(\mathrm{H})_{2^{-}}$ $\left(\mathrm{PH}_{3}\right)_{2}$ are given in Figure 1, along with schematic illustrations of orbitals. All the valence occupied orbitals are shown in the figure, except for PH bonding orbitals which lie in the region of -0.6 to -0.5 hartree.

In the linear $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ complex d orbital levels are split into three groups, two of which are degenerate as shown in Figure 1.

[^3]

Figure 1. MO levels of trans- and cis- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ as well as those of their $\operatorname{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ moieties. The MO symmetry is that of $C_{2 h}$ in the lefthand side and that of $C_{2 v}$ in the right-hand side.
The highest occupied molecular orbital (HOMO) is mainly the $\mathrm{d}_{z^{2}}$ orbital of the Pt atom, as anticipated from the ligand field theoretical consideration. The degenerate lowest unoccupied molecular orbital (LUMO) consists mainly of $P t p_{x}$ and $\mathrm{F}_{y}$ orbitals. Two PtP bonding orbitals lie below d orbital levels, around -0.4 to -0.6 hartree. One of them consists of Pts and $\mathrm{d}_{z^{2}}$ orbitals and $\mathrm{PH}_{3}$ lone-pair orbitals. The other is made of the $\mathrm{Pt} \mathrm{p}_{z}$ orbital and $\mathrm{PH}_{3}$ lone pairs. These bonding features in the linear $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ complex are recognized also in the Mulliken population, as given in the first column of Table I. Gross populations of 0.96 in $s$ valence orbitals and of 0.18 in $p$ valence orbitals of Pt are evidence of participation into the PtP bond formation by these orbitals, whose gross populations are zero in a free $\mathrm{d}^{10}$ singlet Pt atom. The sum of d orbital gross populations is 9.38 , though the two-coordinate $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ complex is formally called a $\mathrm{d}^{10}$ complex.
In forming the trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ complex two hydrogen atoms interact with these MO's to give two bonding orbitals as shown in Figure 1. One is an in-phase mix of the symmetric hydride orbital with $\mathrm{Pt} \mathrm{d}_{y^{2}-z^{2}}$ orbital, and the other consists of the originally vacant $\mathrm{Pt} \mathrm{p}_{y}$ orbital and antisymmetric hydride orbitals. The highest occupied $\mathrm{a}_{\mathrm{g}}$ orbital in $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ becomes unoccupied
a


b




$$
\begin{gathered}
-438081 \\
D_{30}
\end{gathered}
$$

$\mathrm{P}_{1}\left(\mathrm{PH}_{3}\right)_{2}$
$-44.9600$
$\mathrm{C}_{2} \mathrm{~V}$
$\mathrm{C}\left(5-\mathrm{P}_{1}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}\right.$
$-44.9202$

$$
c_{2 v}
$$

TRANSITION STATE

Figure 2. (a) Fully optimized geometries (in $\AA$ and deg ) of $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, and the transition state with their total energies in the smaller basis set. Arrows in the transition state show the reaction coordinate vectors. (b) Fully optimized geometries (in $\AA$ and deg) of $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, and the transition state (TS) with their total energies by using Noell and Hay's basis set and ECP. ${ }^{246}$ The geometry of TS is converged only up to a few thousands of a angstrom, with the energy error of a few tenths of a mircohartree.
in the trans compound and the lowest unoccupied $b_{u}$ orbital becomes occupied as shown in the figure. This switching between an occupied and a vacant orbital suggests that the trans addition of $\mathrm{H}_{2}$ is symmetry forbidden. ${ }^{37}$ The bonding scheme in trans$\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ also can be seen from orbital populations given in Table I. A prominent increase of population in the $\mathrm{Pt} \mathrm{p}_{y}$ orbital and a remarkable decrease in $\mathrm{Pt} y^{2}$ and $z^{2}$ Cartesian d orbitals show the bonding nature described above.

In a bent $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ whose geometry is set to be the same as that of cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, the d levels change their energies from those of the linear structure (right-hand side of Figure 1). The important change is in HOMO; the $\mathrm{d}_{y z}$ orbital becomes the highest occupied. LUMO becomes the hybridized Pt s-p ${ }_{y}$ orbital, directed to the opposite side of $\mathrm{PH}_{3}$ ligands. When two hydrogen atoms are complexed to form cis- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, these HOMO and LUMO interact with antibonding and bonding combinations of H 1s orbitals, respectively. Consequently a decrease in electron density in the $\mathrm{d}_{y z}$ orbital and an increase in the Pt s - p -hybridized orbital take place, as is seen in the gross population in Table I. In contrast to the trans compound, the switching between an occupied and a vacant orbital does not take place in the cis compound.

In transition-metal chemistry the electron count on the hydride is usually presumed to be 2 and molecular structures and reactions are interpreted on the basis of this presumption. However, the gross population asserts almost neutral H ligands in both Pt $(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ complexes; the net charge on H in Table I is between +0.05 and -0.03 , to be compared with Noell and Hay's $+0.02 .{ }^{24 \mathrm{a}}$ The net charge on Pt from the Mulliken population analysis is -0.23 to -0.67 in Table I and -0.45 in Noell and Hay's. These negative charges are in conflict with an intuition for an oxidized Pt compound. This is a defect of the population analysis. In fact Noell and Hay found a positive charge of +1.22 in a modified analysis.

The bonding schemes in cis- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ and trans$\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ are essentially the same as in cis- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and trans- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, respectively, although the molecular symmetry is lower here. Mulliken populations for the hydride methyl compounds are given in the Table I. One notes that the

$\mathrm{Cl} 5-\mathrm{P}_{1}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$


TRANS-P1( H )( $\left.\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$

$-725.9706$
$C_{5}$
transition state

Figure 3. Fully optimized geometries (in $\AA$ and deg) of cis- $\mathrm{Pt}(\mathrm{H})$ $\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, trans- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, and the transition state with their total energies in the smaller basis set. Arrows in the transition state show the reaction coordinate vectors.

Table II. Calculated and Observed PtP Bond Distances ( $\AA$ )

|  | basis |  |
| :---: | :---: | :---: |
|  | smaller | larger ${ }^{\text {a }}$ |
| Calculated |  |  |
| $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ | 2.154 | 2.175 |
| cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ | 2.194 | 2.242 |
| trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ | 2.166 | 2.182 |
| cis- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ | 2.179 |  |
| trans $-\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ | $2.164$ |  |
|  | $2.165$ |  |
| Observed |  |  |
| $\mathrm{Pt}\left[\mathrm{PPh}(t-\mathrm{Bu})_{2}\right]_{2}{ }^{\text {b }}$ | 2.252 |  |
| $\left[\mathrm{Pt}(t-\mathrm{Bu})_{2} \mathrm{P}\left(\mathrm{CH}_{2}\right)_{3} \mathrm{P}(t-\mathrm{Bu})_{2}\right]_{2}{ }^{c}$ | 2.270 |  |
| $\mathrm{PtCl}_{2}\left[(t-\mathrm{Bu})_{2} \mathrm{P}\left(\mathrm{CH}_{2}\right)_{3} \mathrm{P}(t-\mathrm{Bu})_{2}\right]^{¢}$ | 2.281 |  |
| $\mathrm{PtO}_{2}\left[\mathrm{PPh}(t-\mathrm{Bu})_{2}\right]_{2}{ }^{\text {d }}$ | 2.290 |  |
| cis $-\mathrm{PtCl}_{2}\left[\mathrm{P}\left(\mathrm{CH}_{3}\right)_{3}\right]_{2}{ }^{e}$ | 2.239 |  |
| trans $-\mathrm{PtCl}_{2}\left[\mathrm{P}\left(\mathrm{C}_{2} \mathrm{H}_{5}\right)_{3}\right]_{2}{ }^{\text {d }}$ | 2.300 |  |
| trans - $\mathrm{Pt}(\mathrm{H})(\mathrm{Cl})\left[\mathrm{PPh}_{2}\left(\mathrm{C}_{2} \mathrm{H}_{5}\right)\right]_{2}{ }^{\mathrm{g}}$ | 2.268 |  |

${ }^{a}$ Reference 26. ${ }^{b}$ Reference 3. ${ }^{c}$ Reference 5, ${ }^{d}$ Reference 17. ${ }^{e}$ Reference $39 .{ }^{\prime}$ Reference $40 .{ }^{8}$ Reference 41.
$\mathrm{Pt}-\mathrm{C}$ overlap population is smaller than the corresponding $\mathrm{Pt}-\mathrm{H}$ population.
(B) Equilibrium Geometries. The calculated equilibrium geometries are given in Figures 2a and 3. We did not assume any symmetry of the complexes except $C_{s}$, which was necessary for stopping a rather free internal rotation of $\mathrm{PH}_{3}$ groups. The geometry optimization resulted in $D_{2 d}, C_{2 v}$, and $C_{2 h}$ symmetries for $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, and trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, respectively. All the vibrational frequencies except for the near-zero $\mathrm{PH}_{3}$ rotational frequencies are found to be real, confirming that these are in fact equilibrium geometries. The linear structure of $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ is expected from a Walsh diagram for $\mathrm{ML}_{2}$ transi-tion-metal complexes with a $\mathrm{d}^{10}$ electron configuration. ${ }^{38}$ The X -ray structural analysis has actually revealed the PPtP bond angle of $177^{\circ}$ in $\mathrm{Pt}\left[\mathrm{PPh}(t-\mathrm{Bu})_{2}\right]_{2}{ }^{3}$ The square-planar structures of cisand trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ are reasonable, since four-coordinate platinum(II) complexes have been well-known to take a squareplanar structure. ${ }^{6}$
The calculated HPt bond lengths range from 1.519 to 1.563 $\AA$ (Figures 2 a and 3 ), with which there is no experimental structural data to compare. The relativistic ECP calculation on the PtH diatomic molecule gives the bond length of $1.462 \AA,{ }^{24 a}$ whereas it is experimentally known as $1.528 \AA .{ }^{29}$ This suggests that our calculations also underestimate the bond length by about $0.07 \AA$.
The calculated PtP bond distances are compared in Table II with known experimental values. The calculated values of

[^4]Table III. Vibrational Frequencies of $\mathrm{H}_{2}, \mathrm{PH}_{3}, \mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, cis- and trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, and the Transition State ${ }^{a}$

| $\mathrm{H}_{2}{ }^{\text {b }}+\mathrm{PH}_{3}$ | $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ | transitn state ${ }^{c}$ | $\begin{aligned} & \mathrm{Pt}(\mathrm{H})_{2^{-}} \\ & \left(\mathrm{PH}_{3}\right)_{2}{ }^{d} \end{aligned}$ |  | sym and assignt for cis $-\mathrm{Pt}(\mathrm{H})_{2}-$ $\left(\mathrm{PH}_{3}\right)_{2}{ }^{e}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | cis | trans |  |
|  | 160 | 150 (140) | 150 | 190 | $\begin{gathered} \mathrm{A}_{1} \operatorname{PPtPb} \mathrm{~b} \\ \quad \text { (in-plane) } \end{gathered}$ |
|  | 410 | 400 (400) | 380 | 410 | $\mathrm{A}_{1} \mathrm{PtP} \mathrm{s}$ (sym) |
|  | 390 | 440 (410) | 390 | 400 | $\mathrm{B}_{1} \mathrm{PtP} \mathrm{s}$ (anti) |
|  | 180 | 170 (170) | 470 | 210 | $\begin{aligned} & \mathrm{A}_{2} \operatorname{PPtP~b} \\ & \text { (out-of-plane) } \end{aligned}$ |
|  | 560 | 600 (580) | 610 | 630 | $\mathrm{A}_{2} \mathrm{PtPH} b$ (anti out-of-plane) |
|  | 560 | 640 (630) | 640 | 670 | $\mathrm{B}_{1} \mathrm{PtPH} \mathrm{b}$ (anti in-plane) |
|  | 630 | 710 (700) | 730 | 730 | $\begin{aligned} & \mathrm{A}_{1} \mathrm{PtPH} \mathrm{~b} \\ & \text { (sym } \\ & \text { in-plane) } \end{aligned}$ |
|  | 630 | 610 (600) | 760 | 680 | $\mathrm{B}_{2} \mathrm{PtPH} \mathrm{b}$ (sym out-of-plane) |
|  |  | 530 (380) | 920 | 880 | $\mathrm{B}_{2} \mathrm{HPtH} \mathrm{b}$ (out-of-plane) |
|  |  |  | 920 | 920 | $\mathrm{B}_{1} \mathrm{HPtP} \mathrm{b}$ |
|  |  | 330 (250) | 1160 | 1090 | $\mathrm{A}_{1} \mathrm{HPtH} \mathrm{b}$ <br> (in-plane) |
| 1430 [992] | 1240 | 1280 (1280) | 1270 | 1260 | $\mathrm{B}_{1} \mathrm{HPH} \mathrm{b}\left(\mathrm{a}_{1}\right)$ |
|  | 1260 | 1310 (1310) | 1300 | 1260 | $\mathrm{A}_{1} \mathrm{HPH} \mathrm{b}\left(\mathrm{a}_{1}\right)$ |
| 1540 [1118] | 1400 | 1420 (1420) | 1420 | 1390 | $\mathrm{B}_{1}$ HPH b (e) |
|  | 1400 | 1430 (1430) | 1420 | 1390 | $\mathrm{A}_{2}$ HPH b (e) |
| 1540 [1118] | 1410 | 1440 (1440) | 1420 | 1400 | $\mathrm{A}_{1}$ HPH b (e) |
|  | 1410 | 1440 (1440) | 1420 | 1400 | $\mathrm{B}_{2} \mathrm{HPH} \mathrm{b}$ (e) |
|  |  |  | 2140 | 1810 | $\mathrm{B}_{1} \mathrm{PtHss}$ (anti) |
|  |  | 360i (270i) | 2210 | 2150 | $\mathrm{A}_{1} \mathrm{PtHss}$ (sym) |
| 3100 [2323] | 3050 | 3050 (3050) | 3060 | 3070 | $\mathrm{A}_{1} \mathrm{PH} s\left(\mathrm{a}_{1}\right)$ |
|  | 3050 | 3060 (3060) | 3060 | 3070 | $\mathrm{B}_{1} \mathrm{PH} \mathrm{s}\left(\mathrm{a}_{1}\right)$ |
| 3150 [2328] | 3080 | 3080 (3080) | 3090 | 3100 | $\mathrm{A}_{2} \mathrm{PH} s$ (e) |
|  | 3080 | 3080 (3080) | 3090 | 3100 | $\mathrm{B}_{2} \mathrm{PH} s$ (e) |
| 3150 [2328] | 3090 | 3100 (3100) | 3120 | 3110 | $\mathrm{A}_{1} \mathrm{PH} \mathrm{s}$ (e) |
|  | 3100 | 3100 (3100) | 3120 | 3120 | $\mathrm{B}_{1} \mathrm{PH} \mathrm{s}$ (e) |
| 4660 (3300) |  | 4150 (2940) |  |  | $\mathrm{A}_{1} \mathrm{HH} \mathrm{s}$ |

${ }^{a}$ The force constant matrix is obtained at the smaller basis RHF level. Vibrational modes corresponding to internal rotation of $\mathrm{PH}_{3}$ are excluded. ${ }^{b}$ Values in square brackets are experimental fundamental frequencies for gaseous $\mathrm{PH}_{3} .{ }^{42}$ A values in parentheses is a calculated frequency of $\mathrm{D}_{2}$. ${ }^{c}$ Vibrational frequencies of $\mathrm{Pt}(\mathrm{D})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ are given in parentheses. Three near-zero frequencies are not given. Two of these correspond to $\mathrm{PH}_{3}$ rotation, and the third corresponds to $\mathrm{H}_{2}$ out-ofplane translational mode. ${ }^{d}$ Two near-zero frequencies corresponding to $\mathrm{PH}_{3}$ rotations are not given. ${ }^{e}$ The b and s means bending and stretching, respectively. For modes within $\mathrm{PH}_{3}$ ligands irreducible representation in $C_{3 v}$ symmetry is given in parentheses.
$2.15-2.20 \AA$, regardless whether we use a smaller (STO-2G) or a larger (LP-31G) basis set for P, are about $0.1 \AA$ shorter than 2.2-2.3 $\AA$, the observed values for many Pt phosphine compounds. ${ }^{6}$ This is probably due to the inadequancy of the Pt ECP parameters adopted in the calculation. The poor basis set used for $P$ and the lack of correlation effect may also be responsible for this error.

Noell and Hay reported optimized HPt distance and HPtH and PPtP bond angles of cis- and trans- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ by the ener-gy-only method with a fixed PtP bond distance of $2.268 \AA .{ }^{246}$ Their HPt bond distances of 1.55 and $1.61 \AA$ in the cis and trans compounds, respectively, are about $0.04 \AA$ longer than those of the present work. Concerned with this difference, we have carried out a full geometry optimization for cis- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ by using the ECP parameters and the basis set used by them. ${ }^{24 b}$ The result, shown in the second column of Figure 2 b , gives the HPt distance of $1.51 \AA$, in agreement with the present work, and the PPt distance of $2.51 \AA$, much longer than the present results and those observed in experiments.

The HPt bond length becomes shorter as the trans-positioned ligand is switched from $\mathrm{CH}_{3}$ to H and from H to $\mathrm{PH}_{3}$. The same trend is found for the PPt and the CPt bond lengths. From this fact, one can deduce the trans influence order $\mathrm{H}^{-}>\mathrm{CH}_{3}{ }^{-}>\mathrm{PH}_{3}$. The trans influence order will also be discussed in section E .

Table IV. Vibrational Frequencies of $\mathrm{CH}_{4}$, cis- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, trans- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, and the Transition State ${ }^{a}$

| $\mathrm{CH}_{4}{ }^{\text {b }}$ | $\underline{\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}}$ |  | transitn state ${ }^{c}$ | $\begin{gathered} \text { symm and } \\ \text { assignt for } \\ \text { cis-Pt(H)- } \\ \left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}{ }^{d} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: |
|  | cis ${ }^{\text {c }}$ | trans |  |  |
|  | 160 (160) | 120 | 120 (120) | $\begin{gathered} \mathrm{A}^{\prime \prime} \mathrm{PPtC} \mathrm{~b} \\ \text { (out-of- } \\ \text { plane) } \end{gathered}$ |
|  | 170 (170) | 160 | 100 (100) | $\mathrm{A}^{\prime} \operatorname{PPtPb}$ (in-plane) |
|  | 250 (250) | 230 | 150 (150) | $\mathrm{A}^{\prime} \mathrm{PPtC}$ b (in-plane) |
|  | 370 (370) | 400 | 370 (370) | $\mathrm{A}^{\prime} \mathrm{PtP} \mathrm{s}$ (sym) |
|  | 400 (400) | 400 | 400 (400) | $\mathrm{A}^{\prime} \mathrm{PtP} \mathrm{s}$ (anti) |
|  | 490 (410) | 600 | 430 (350) | $\mathrm{A}^{\prime \prime} \mathrm{PPtH} \mathrm{b}$ (out-ofplane) |
|  | 560 (560) | 520 | 520 (520) | $\mathrm{A}^{\prime} \mathrm{PtC} \mathrm{s}$ |
|  | 640 (640) | 630 | 580 (570) | $\begin{aligned} & \mathrm{A}^{\prime \prime} \mathrm{PtPH} \mathrm{~b} \\ & \text { (anti } \\ & \text { out-of-plane) } \end{aligned}$ |
|  | 680 (670) | 690 | 640 (640) | $\mathrm{A}^{\prime} \operatorname{PtPH} \mathrm{b}$ (anti in-plane) |
|  | 720 (720) | 740 | 720 (720) | $\mathrm{A}^{\prime}$ PtPH b (sym in-plane) |
|  | 780 (680) | 690 | 620 (610) | $\mathrm{A}^{\prime \prime} \mathrm{PtPH} \mathrm{b}$ (sym out-of-plane) |
|  | 870 (870) | 900 | 1080 (970) | $\mathrm{A}^{\prime \prime} \mathrm{PtCH} \mathrm{b}$ (out-ofplane) |
|  | 940 (940) | 880 | 1090 (980) | $\mathrm{A}^{\prime} \mathrm{PtCH}$ b (in-plane) |
|  | 1150 (840) | 1070 | 1030i (810i) | $\mathrm{A}^{\prime} \mathrm{HPtC}$ b (in-plane) |
|  | 1260 (1260) | 1240 | 1290 (1290) | $\mathrm{A}^{\prime} \mathrm{HPH} \mathrm{b}\left(\mathrm{a}_{1}\right)$ |
|  | 1290 (1290) | 1260 | 1320 (1320) | $A^{\prime}$ HPH b $\left(a_{1}\right)$ |
|  | 1410 (1410) | 1390 | 1420 (1420) | $A^{\prime \prime}$ HPH b (e) |
|  | 1410 (1410) | 1390 | 1430 (1430) | $A^{\prime}$ HPH b (e) |
|  | 1410 (1410) | 1400 | 1430 (1430) | $\mathrm{A}^{\prime} \mathrm{HPH} \mathrm{b}$ (e) |
|  | 1410 (1410) | 1400 | 1440 (1440) | $A^{\prime \prime} \mathrm{HPH} \mathrm{b}$ (e) |
| 1520 (1150) $\mathrm{T}_{2}$ | 1440 (1440) | 1450 | 1450 (1450) | $\mathrm{A}^{\prime} \mathrm{HCH} \mathrm{b}\left(\mathrm{a}_{6}\right)$ |
| 1740 (1230) E | 1660 (1660) | 1670 | 1630 (1530) | $\mathrm{A}^{\prime} \mathrm{HCH} \mathrm{b} \mathrm{(e)}$ |
|  | 1690 (1690) | 1680 | 1680 (1680) | $\mathrm{A}^{\prime \prime} \mathrm{HCH}$ b (e) |
|  | 2160 (1530) | 2030 | 2170 (1710) | $\mathrm{A}^{\prime} \mathrm{PtHs}$ |
|  | 3040 (3040) | 3060 | 3120 (3120) | $\mathrm{A}^{\prime} \mathrm{PH} s\left(a_{1}\right)$ |
|  | 3060 (3060) | 3060 | 3030 (3030) | $\mathrm{A}^{\prime} \mathrm{PH} \mathrm{s}\left(\mathrm{a}_{1}\right)$ |
| 3190 (2250) $\mathrm{A}_{1}$ | 3070 (3070) | 3080 | 3060 (3060) | $\mathrm{A}^{\prime} \mathrm{CH} s\left(\mathrm{a}_{1}\right)$ |
|  | 3070 (3070) | 3090 | 3060 (3060) | $\mathrm{A}^{\prime \prime} \mathrm{PH} \mathrm{s}$ (e) |
|  | 3090 (3090) | 3090 | 3070 (3070) | $\mathrm{A}^{\prime \prime} \mathrm{PH} \mathrm{s}(\mathrm{e})$ |
| 3280 (2430) $\mathrm{T}_{2}$ | 3090 (3090) | 3100 | 3090 (3090) | $\mathrm{A}^{\prime \prime} \mathrm{CH}$ s (e) |
|  | 3100 (3100) | 3100 | 3160 (3160) | $\mathrm{A}^{\prime} \mathrm{PH} \mathrm{s}$ (e) |
|  | 3110 (3110) | 3110 | 3240 (3240) | $\mathrm{A}^{\prime} \mathrm{PH} \mathrm{s}$ (e) |
|  | 3240 (3240) | 3120 | 3290 (3290) | $\mathrm{A}^{\prime} \mathrm{CH} s$ (e) |

${ }^{a}$ Based on the force constants obtained at the smaller basis RHF level. Three near-zero frequencies corresponding to internal rotations of $\mathrm{PH}_{3}$ and $\mathrm{CH}_{3}$ are excluded, ${ }^{b}$ Values in parentheses are calculated $\mathrm{CD}_{4}$ frequencies. The irreducible representations are those in $T_{d}$ symmetry. ${ }^{c}$ Vibrational frequencies of $\mathrm{Pt}(\mathrm{D})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ are given in parentheses. ${ }^{d}$ The b and s means bending and stretching, respectively, For modes within $\mathrm{PH}_{3}$ and $\mathrm{CH}_{3}$ ligands irreducible representation in $C_{30}$ is given in parentheses.

The HPtH bond angle in the cis dihydride complex is calculated to be $81.8^{\circ}$, smaller than $90^{\circ}$, while the $\operatorname{PPtP}$ angle is $104.4^{\circ}$, larger than $90^{\circ}$. This is in agreement with the Hückel calculations for the $\mathrm{R}_{2} \mathrm{ML}_{2}$ transition-metal complexes. ${ }^{19}$ The ligand-ligand steric repulsion is expected to be the smallest in case of $\mathrm{PH}_{3}$ among substituted bulky phosphines. Therefore, the calculated bond angle is considered to be electronically optimum for phosphines that have the same basicity as $\mathrm{PH}_{3}$.
(C) Vibrational Frequencies. All the vibrational modes of $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, cis $-\mathrm{Pt}(\mathrm{H})$ $\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, and trans $-\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ are obtained from force constant matrices calculated by a numerical differentiation

Table V. Calculated and Observed PtH Stretching Frequencies $\left(\mathrm{cm}^{-1}\right)^{a}$

|  | anti sym | sym |
| :---: | :---: | :---: |
| Calculated |  |  |
| trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ | 1810 (1290 ${ }^{\text {d }}$ ) | 2150 (1520 ${ }^{\text {d }}$ ) |
| cis- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ | 2140 (1520 ${ }^{\text {d }}$ ) | 2210 (1570 ${ }^{\text {d }}$ ) |
| trans $-\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ | 2030 |  |
| cis $-\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ | 2160 |  |
| Observed |  |  |
| trans $-\mathrm{Pt}(\mathrm{H})_{2}\left[\mathrm{P}(i-\mathrm{Pr})_{3}\right]_{2}{ }^{\text {b }}$ | 1735 |  |
| trans $-\mathrm{Pt}(\mathrm{H})_{2}\left[\mathrm{P}\left(\mathrm{c}-\mathrm{C}_{6} \mathrm{H}_{41}\right)_{3}\right]_{2}{ }^{\text {b }}$ | 1710 (1230 ${ }^{\text {d }}$ ) |  |
| trans $-\mathrm{Pt}(\mathrm{H})_{2}\left[\mathrm{PPH}(t-\mathrm{Bu})_{2}\right]_{2}{ }^{b}$ | 1780 |  |
| trans $-\mathrm{Pt}(\mathrm{H})_{2}\left[\mathrm{P}\left(\mathrm{C}_{2} \mathrm{H}_{5}\right)_{3}\right]_{2}{ }^{\text {b }}$ | 1725 |  |
| cis $-\mathrm{Pt}(\mathrm{H})_{2}\left[(t-\mathrm{Bu})_{2} \mathrm{P}\left(\mathrm{CH}_{2}\right)_{3} \mathrm{P}(t-\mathrm{Bu})_{2}\right]^{c}$ | 1995 (1430d) |  |
| cis $-\mathrm{Pt}(\mathrm{H})_{2}\left[(t-\mathrm{Bu})_{2} \mathrm{P}\left(\mathrm{CH}_{2}\right)_{2} \mathrm{P}(t-\mathrm{Bu})_{2}\right]^{c}$ | 1980, 2008 |  |
| cis $-\mathrm{Pt}(\mathrm{H})_{2}\left[\mathrm{Men}_{2} \mathrm{P}\left(\mathrm{CH}_{2}\right)_{2} \mathrm{PMen}_{2}\right]^{c}$ | 1989 |  |
| $\begin{aligned} & c i s-\mathrm{Pt}(\mathrm{H})_{2}\left[\mathrm{Men}(\mathrm{Ph}) \mathrm{P}\left(\mathrm{CH}_{2}\right)_{2} \mathrm{P}(\mathrm{Ph})-\right. \\ & \mathrm{Men}]^{c} \end{aligned}$ | 1977, 2004 |  |
| $\begin{gathered} c i s-\mathrm{Pt}(\mathrm{H})_{2}\left[t-\mathrm{Bu}(\mathrm{Ph}) \mathrm{P}\left(\mathrm{CH}_{2}\right)_{2} \mathrm{P}(\mathrm{Ph}) t-\right. \\ \mathrm{Bu}]^{\mathrm{c}} \end{gathered}$ | 1985 |  |
| cis $-\mathrm{Pt}(\mathrm{H})_{2}\left[\mathrm{Ph}_{2} \mathrm{P}\left(\mathrm{CH}_{2}\right)_{2} \mathrm{PPh}_{2}\right]^{\text {c }}$ | 1930 |  |

${ }^{a}$ At RHF level with the smaller basis set. ${ }^{b}$ Reference $4 .{ }^{c}$ Reference 5. ${ }^{d}$ In a deuterated compound $\operatorname{Pt}(\mathrm{D})_{2} \mathrm{~L}_{2}$.

Table VI. Calculated and Observed PtC Stretching Frequencies $\left(\mathrm{cm}^{-1}\right)^{a}$

| Calculated |  |
| :---: | :---: |
| trans $-\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ |  |
| cis- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ |  |
| Observed |  |
| trans- $\mathrm{Pt}(\mathrm{X})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}{ }^{b}$ | 520 |
| $\mathrm{X}=\mathrm{NO}_{3}$ | 560 |
| $\mathrm{X}=\mathrm{NC}$ |  |
| $\mathrm{X}=\mathrm{Cl}$ | 566 |
| $\mathrm{X}=\mathrm{Br}$ | 556 |
| $\mathrm{X}=\mathrm{NO}_{2}$ | 551 |
| $\mathrm{X}=\mathrm{I}$ | 548 |
| $\mathrm{X}=\mathrm{CN}$ | 544 |
| cis- $\mathrm{Ft}(\mathrm{Cl})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}{ }^{b}$ | 540 |
|  | 516 |
|  | 516 |

${ }^{a}$ See the footnote $a$ of Table V. ${ }^{b}$ Reference 43.
of the analytical energy gradients. Calculated normal vibrational frequencies and their assignments, excluding near-zero frequencies of $\mathrm{CH}_{3}$ and $\mathrm{PH}_{3}$ rotations, are shown in Tables III and IV. Experimentally observed PtH and PtC stretching fundamental frequencies are compared in Tables V and VI, respectively, with the calculated normal-mode frequencies.

The agreement is generally very good for both PtH and PtC stretchings with the calculated values overestimating the observed values by $4-8 \%$. Both experimentally and theoretically, the antisymmetric PtHl stretching for the cis dihydride compound is found to have a higher frequency than the trans compound by over $200 \mathrm{~cm}^{-1}$. For the hydride methyl compound the calculation similarly predicts a higher PtC stretching frequency for cis than for trans. These results can be interpreted as an indication that the hydride shows a stronger trans influence than $\mathrm{PH}_{3}$. In the experiment for $\mathrm{Pt}\left(\mathrm{CH}_{3}\right) \mathrm{Cl}\left(\mathrm{PEt}_{3}\right)_{2}$, trans has a higher PtC frequency than cis, reflecting a weaker trans influence of Cl .

Vibrational modes originating from $\mathrm{PH}_{3}$ ligands, namely, HPH bending (calculated $1240-1540 \mathrm{~cm}^{-1}$ ) and PH stretching (calculated $3050-3150 \mathrm{~cm}^{-1}$ ) are up to $40 \%$ too large, as expected from this small STO-2G basis set for P. However, one can clearly recognize a decrease of these frequencies upon coordination to the Pt atom and a subsequent slight reversal upon oxidative addition.

There are many vibrational modes whose frequencies are about $1000 \mathrm{~cm}^{-1}$ or less. These frequencies have not yet been identified experimentally. The present calculation predicts about 900 and $1000-1200 \mathrm{~cm}^{-1}$ for the HPtH out-of-plane and in-plane motions, respectively. However, it should be remembered that the present

[^5]Table VII. Calculated Relative Energy for $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}+\mathrm{H}_{2} \rightarrow$ $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}(\text { in } \mathrm{kcal} / \mathrm{mol})^{a}$

| method | $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ <br> $+\mathrm{H}_{2}{ }^{b}$ | transition <br> state | $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ |  |
| :--- | :---: | :---: | :---: | :---: |
|  | 0 | +5.2 | -36.9 | cis |
| RHF | 0 | +88.7 | -27.0 | -25.1 |
| SD-CI | 0 | +7.3 | -26.8 | -24.0 |
| SD-CI + QC | 0 | +8.2 | -21.5 | -20.3 |
| SD-CI + QC + ZPC | 0 | +8.2 |  |  |

${ }^{a}$ The larger basis set. ${ }^{b} \mathrm{RHF}$ energy for $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}=-43.61521$ hartree and for $\mathrm{H}_{2}=-1.12295$ hartree. $\mathrm{SD}-\mathrm{CI}$ correlation energy $=$ -0.29819 and -0.02476 hartree. $\mathrm{QC}=-0.04193$ and -0.00034 hartree. $\mathrm{ZPC}=43.2$ and $6.7 \mathrm{kcal} / \mathrm{mol}$.

Table VIII. Calculated Relative Energy for $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}+\mathrm{CH}_{4} \rightarrow$ $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}(\text { in } \mathrm{kcal} / \mathrm{mol})^{a}$

| method | $\begin{gathered} \mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2} \\ +\mathrm{CH}_{4}^{b} \end{gathered}$ | transitn state | $\begin{gathered} \mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)- \\ \left(\mathrm{PH}_{3}\right)_{2} \\ \hline \end{gathered}$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  |  | cis | trans |
| RHF | 0 | 22.6 | -15.0 | -17.5 |
| SD-CI | 0 | 35.3 | 1.2 | 1.2 |
| $\mathrm{SD}-\mathrm{CI}+\mathrm{QC}$ | 0 | 29.6 | -4.7 | -3.9 |
| $\mathrm{SD}-\mathrm{CI}+\mathrm{QC}+\mathrm{ZPC}$ | 0 | 28.2 | -4.5 | -3.5 |

${ }^{a}$ The larger basis set. ${ }^{b}$ RHF energy for $\mathrm{CH}_{4}=-39.97688$ hartree. $\mathrm{SD}-\mathrm{CI}$ correlation energy $=-0.11268$ hartree. $\mathrm{QC}=-0.00524$ hartree. $\mathrm{ZPC}=30.1 \mathrm{kcal} / \mathrm{mol}$.
level of calculations is expected to overestimate experimental frequencies by $5-15 \%{ }^{44}$ Both of these frequencies in cis-Pt$(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ are larger than those in trans- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$. Bending and stretching modes involving a PtP bond occur at $700 \mathrm{~cm}^{-1}$ or less. These frequencies are less accurate than those involving a PtH bond, because of the minimal (STO-2G) $\mathrm{PH}_{3}$ basis set.
(D) Relative Stability of Four-Coordinate Cis and Trans Platinum(II) Compounds. In the oxidative addition reaction of $\mathrm{H}_{2}$ to platinum(0) diphosphine only the trans isomer has been found as a product. The relative stability of cis and trans isomers could be a determining factor for this selectivity. ${ }^{45}$ We have calculated the energy of the two structures relative to the reactant by the RHF and the SD-CI + QC method with a larger basis set for the RHF-optimized geometries. CI calculations give the coefficient of the RHF configuration of about 0.924-0.930 for all the Pt complexes studied, supporting qualitative discussions on electronic structures based on an RHF wave function. The zero-point energy correction (ZPC) has been calculated with the vibrational frequencies in Tables III and IV and therefore is probably overestimated. The results with ZPC are shown in the last two columns of Tables VII and VIII for $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, respectively.
In our best calculation ( $\mathrm{SD}-\mathrm{CI}+\mathrm{QC}+\mathrm{ZPC}$ ) the energy differences between cis and trans compounds both in Tables VII and VIII are less than a few kilocalories per mole, which is within the error limit of the present level of calculation. This means that no obvious thermodynamical preference of the trans compound has been found in our model compound $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$. The trans selectivity found experimentally will be explained later in terms of a steric repulsion between bulky phosphines.
The difference in the principal mode of reaction, i.e., the ox-idative-addition reaction for $\mathrm{H}_{2}$ vs. the reductive-elimination reaction for $\mathrm{CH}_{4}$, can be understood easily in terms of the difference in exothermicity, calculated in Tables VII and VIII. For $\mathrm{H}_{2}$, the addition product is more stable than the two-coordinate reactant by about $20 \mathrm{kcal} / \mathrm{mol}$, making the addition reaction more favorable. For $\mathrm{CH}_{4}$, the reaction is found to be essentially thermoneutral, and the entropy factor should favor the elimination product. As discussed in the following, the exact exothermicity is hard to evaluate. However, the qualitative difference in behavior between $\mathrm{H}_{2}$ and $\mathrm{CH}_{4}$ should be meaningful.

[^6]Table IX. Bond Strength, Bond Distance, Stretching Frequency, and Overlap Population of PtH, PtC , and PtP

|  | $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ | $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ |  | $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\left(\mathrm{PH}_{3}\right)_{2}\right.$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | cis | trans | cis | trans |
| Bond Strength (kcal/mol) |  |  |  |  |  |
| PtH |  | $70^{a}$ | $61^{a}$ | $(70)^{d}$ | (65) ${ }^{\text {d }}$ |
| PtC |  |  |  | (34) ${ }^{\text {e }}$ | $(41)^{e}$ |
| PtP | $40^{\text {b }}$ | $38^{\text {c }}$ | $45^{\text {c }}$ |  |  |
| Bond Distance ( $\AA$ ) |  |  |  |  |  |
| PtH |  | 1.522 | 1.563 | 1.519 | 1.552 |
| PtC |  |  |  |  |  |
| PtP | 2.154 | 2.194 | 2.166 | 2.179, 2.200 | $2.164,2.165$ |
| Stretching ( $\mathrm{cm}^{-1}$ ) |  |  |  |  |  |
| PtH |  | 2140, 2210 | 1810, 2150 | 2160 | 2030 |
| PtC |  |  |  | 560 | 520 |
| PtP | 390, 400 | 380,390 | 400, 410 | 370, 400 | 400 |
| Bond Population |  |  |  |  |  |
| PtH |  | 0.39 | 0.38 | 0.39 | 0.38 |
| PtC |  |  |  | 0.32 | 0.29 |
| PtP | 0.35 | 0.26 | 0.33 | 0.25 | 0.33 |
| ${ }^{a}\left[E\left(\mathrm{Pt}_{( }\left(\mathrm{PH}_{3}\right)_{2}\right)+2 E(\mathrm{H})-E\left(\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}\right)\right] / 2 .{ }^{b} E\left(\mathrm{PtPH}_{3}\right)+E$ -$\left(\mathrm{PH}_{3}\right)-E\left(\mathrm{Pt}_{( }\left(\mathrm{PH}_{3}\right)_{2}\right) \cdot{ }^{c} E\left(\mathrm{Pt}(\mathrm{H})_{2} \mathrm{PH}_{3}\right)+E\left(\mathrm{PH}_{3}\right)-E\left(\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}\right)$. ${ }^{d}$ Assumed value. See Text. ${ }^{e} E\left(\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}\right)+E(\mathrm{H})+E\left(\mathrm{CH}_{3}\right)-E$ -$\left(\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}\right)-E_{\mathrm{A}}(\mathrm{PtH})$, where $E_{\mathrm{A}}(\mathrm{PtH})$ is the assumed value of PtH bond strength. |  |  |  |  |  |

Noell and Hay calculated the stabilization energies ${ }^{24 \mathrm{a}}$ of cis$\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and trans $-\mathrm{Pt}\left(\mathrm{H}_{2}\left(\mathrm{PH}_{3}\right)_{2}\right.$ relative to the $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ $+\mathrm{H}_{2}$ system and obtained from the RHF calculation the energies of 6.7 and $10.6 \mathrm{kcal} / \mathrm{mol}$, respectively, which are about 30 $\mathrm{kcal} / \mathrm{mol}$ less than those of 36.9 and $38.0 \mathrm{kcal} / \mathrm{mol}$ at the RHF level of calculation in the present work. The dissociation energy, $D_{\mathrm{e}}$, of the ground state of the PtH molecule is to be $53.4 \mathrm{kcal} / \mathrm{mol}$ in their basis set and ECP ${ }^{24 \mathrm{a}}$ while to be $65.0 \mathrm{kcal} / \mathrm{mol}$ in the present set and ECP. The difference of $12 \mathrm{kcal} / \mathrm{mol}$ per each PtH bond will explain a main part of the stabilization energy difference. Because the dissociation energy of the PtH molecule has not been known experimentally, we cannot determine which result is more reasonable. It is noted, however, that the calculated vibrational frequency of PtH molecule is $2260 \mathrm{~cm}^{-1}$ in the present calculation and is in better agreement with the observed value of $2377 \mathrm{~cm}^{-1}$ than that of $2009 \mathrm{~cm}^{-1}$ by Noell and Hay. ${ }^{24 a}$
(E) Nature of PtH, PtC, and PtP Bonds. A comparison between cis and trans conformers of $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{Pt}(\mathrm{H})(\mathrm{C}$ $\left.\mathrm{H}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ gives interesting features of $\mathrm{PtH}, \mathrm{PtC}$, and PtP bonds. In Table IX, we have compared the bonding characteristics such as estimated bond strength, bond distance, stretching frequency, and bond population among the compounds we studied. The PtH bond strength is calculated to be half the energy difference between $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}+\mathrm{H}+\mathrm{H}$ and $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and the PtP bond strength to be an energy difference between $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)+\mathrm{PH}_{3}$ and $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$. The calculated PtH bond strength is 70 and 61 $\mathrm{kcal} / \mathrm{mol}$ for cis- and trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$, respectively. Even though a direct comparison with experiments cannot be made, it is interesting to note that Beauchamp et al. and Ridge et al. estimated the metal cation-hydride bond energy to be in the range of $50-80 \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{Cr}, \mathrm{Mn}, \mathrm{Fe}, \mathrm{Co}, \mathrm{Ru}$, and Rh complexes. ${ }^{46.47}$ Halpern also estimated a bond dissociation energy of MH as $\sim 60$ $\mathrm{kcal} / \mathrm{mol}^{48}$ and that of MC as $20-30 \mathrm{kcal} / \mathrm{mol}^{49}$ We find that the PtH bond in the cis compound has a larger bond strength,

[^7]a shorter distance, a larger vibrational frequency, and a larger bond population than those in the trans compound. To the contrary, the PtP bond in the cis product has a smaller bond strength, a longer distance, lower vibrational frequencies, and a smaller bond population. The PtH bond in cis appears to be stronger than in trans, whereas the PtP bond in cis is weaker. These results can be explained by a difference in the so-called trans influence; the trans influence of H is stronger than that of $\mathrm{PH}_{3}$. This is also consistent with a longer PtCl bond distance of $2.42 \AA$ in trans$\mathrm{Pt}(\mathrm{H})(\mathrm{Cl})\left[\mathrm{PPh}_{2}\left(\mathrm{C}_{2} \mathrm{H}_{5}\right)\right]_{2},{ }^{39}$ where H is in the trans position to Cl , than that of $2.37 \AA$ in cis $-\mathrm{PtCl}_{2}\left[\mathrm{P}\left(\mathrm{CH}_{3}\right)_{2}\right]_{2},{ }^{41}$ where phosphine is in the trans position to Cl .

In cis- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, both of PtH and PtC bonds have smaller bond distances, larger stretching frequencies, and larger bond populations than those in the corresponding trans compound. These results suggest that the trans influence of a $\mathrm{CH}_{3}$ ligand is larger than that of $\mathrm{PH}_{3}$ ligand. In the trans dihydride the PtH bond distance is larger than that in trans hydride methyl compound, indicating a smaller trans influence of $\mathrm{CH}_{3}$ than H .
In order to estimate the PtC bond strength, we first calculated the stabilization energy of cis- and trans- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ relative to $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}+\mathrm{H}+\mathrm{CH}_{3}$ to be 104 and $106 \mathrm{kcal} / \mathrm{mol}$, respectively. If we assume and subtract the PtH bond strength of 70 and $65 \mathrm{kcal} / \mathrm{mol}$, respectively, we obtain an estimate of the PtC bond strength to be 34 and $41 \mathrm{kcal} / \mathrm{mol}$ for cis- and trans$\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, respectively. The assumed PtH bond length of the cis compound is that of cis- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and is probably reasonable because in both of them the ligand trans to hydride is a phosphine. The assumed PtH bond strength of $65 \mathrm{kcal} / \mathrm{mol}$ for trans $-\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ is an average of the PtH bond strengths in cis- and trans $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and is deduced from the fact that a ligand trans to hydride is $\mathrm{CH}_{3}$ in trans $-\mathrm{Pt}(\mathrm{H})$ $\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ and that a trans influence of $\mathrm{CH}_{3}$ group is weaker than hydride and stronger than phosphine. The estimated PtC bond strengths of 34 and $41 \mathrm{kcal} / \mathrm{mol}$ are in good agreement with a range of values of $20-30 \mathrm{kcal} / \mathrm{mol}^{49}$ for metal-alkyl bond dissociation energies of coordinately saturated cobalt compounds and a value of $41 \mathrm{kcal} / \mathrm{mol}^{46}$ for the $\mathrm{Co}-\mathrm{CH}_{3}$ bond obtained from ion-beam studies.
Transition States for Oxidative Addition of $\mathrm{H}_{2}$ to $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ and Reductive Elimination of $\mathrm{CH}_{4}$ from $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$. (A) Transition State for $\mathrm{H}_{2}+\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2} \rightarrow \mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$. We have found a transition state for the reaction $\mathrm{H}_{2}+\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2} \rightarrow \mathrm{Pt}$ $(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and shown its geometry in the last column in Figure 2a. In the search of the transition state $C_{s}$ symmetry is assumed to restrict rather free rotations of $\mathrm{PH}_{3}$ around metal-phosphorus bonds. The resultant transition state has $C_{20}$ symmetry. Nor-mal-mode frequencies of the transition state are given in Table III where near-zero vibrational frequencies are not included.
Relative energies of the transition state in various approximations are given in Table VII. Our best calculation ( $\mathrm{CI}+\mathrm{QC}$ ) gives a classical barrier height of $7 \mathrm{kcal} / \mathrm{mol}$. The zero-point energy correction (ZPC) changes the effective barrier to 8 $\mathrm{kcal} / \mathrm{mol}$. Though these values should be taken to be only semiquantitative, it is safe to say that the barrier for this model reaction is low, consistent with the experimental fact ${ }^{4.5}$ that the oxidative addition of $\mathrm{H}_{2}$ takes place easily.

The geometry of the transition state resembles that of the reactants; the HH bond length of $0.77 \AA$ is only $4 \%$ longer than $0.73 \AA$ in the $\mathrm{H}_{2}$ molecule, the PPtP angle of $148^{\circ}$ is not far from $180^{\circ}$ of the reactant $\operatorname{Pt}\left(\mathrm{PH}_{3}\right)_{2}$. On the other hand, the newly formed PtH distance is still substantially longer than that in the product. The transition state for the oxidative addition is, therefore, characterized as "an early transition state".
In a kinetic study of hydrogen and deuterium addition reaction of $\operatorname{IrCl}(\mathrm{CO})\left(\mathrm{PPh}_{3}\right)_{2}$, it has been suggested ${ }^{50}$ on the basis of a small kinetic isotope effect ${ }^{9}$ that the HH bond stretch is small at the transition state. The present result is consistent with this suggestion. Dedieu and Strich also suggested small HH stretching in the transition state based on their extended Hückel study of

[^8]Table X. Analysis of Kinetic Deuterium Isotope Effect at 300 K for the Oxidative-Addition Reaction $(\mathrm{X})_{2}+\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2} \rightarrow \mathrm{Pt}(\mathrm{X})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ $(\mathrm{X}=\mathrm{H} \text { and } \mathrm{D})^{a}$

| sym | reactant |  | transitn state |  | $R_{i}{ }^{*} R_{i}^{-1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\begin{aligned} & \hline \text { assignt }{ }^{b} \\ & \text { or } \nu, \mathrm{cm}^{-1} \end{aligned}$ | $R_{i}$ |  |  |  |
|  |  |  | $\nu,{ }^{\text {b }} \mathrm{cm}^{-1}$ | $R_{i}^{*}$ |  |
| $\mathrm{A}_{1}$ | $t(x)$ | 1.0 | 360 i | 1.31 | 1.31 |
|  | 610 | 1.0 | 710 | 0.99 | 0.99 |
|  | 4660 | 0.054 | 4160 | 0.077 | 1.43 |
| $\mathrm{A}_{2}$ | $r(x)$ | 1.0 | 530 | 0.89 | 0.89 |
|  | 580 | 1.0 | 600 | 0.98 | 0.98 |
| $\mathrm{B}_{1}$ | $t(z)$ | 1.0 | 330 | 0.96 | 0.96 |
|  | $r(y)$ | 1.0 | 440 | 0.98 | 0.98 |
| $\mathrm{B}_{2}$ | $t(y)$ | 1.0 |  | $(1.0)^{\text {c }}$ | $(1.0)^{\text {c }}$ |
| total ${ }^{\text {d }}$ |  | 0.054 |  | 0.061 | 1.48 |

${ }^{a}$ Only factors contributing significantly to the isotope effect are listed. ${ }^{b}$ The $t(i)$ is a translation of $\mathrm{X}_{2}$ to the $i$ direction, and the $r(i)$ is the rotation of $\mathrm{X}_{2}$ around $i$ axis based on the coordinate system in Scheme I. Frequencies are those in dihydride compound $\mathrm{Pt}(\mathrm{H})_{2}(\mathrm{P}-$ $\left.\mathrm{H}_{3}\right)_{2}$. ${ }^{c}$ Assumed value. See text. ${ }^{d}$ Total product of contributions of all the modes where $R_{i}$ and $R_{i}{ }^{*}$ for $\mathrm{PH}_{3}$ rotations having near-zero frequencies are assumed to be 1 .
$\mathrm{RhCl}\left(\mathrm{PH}_{3}\right)_{3}+\mathrm{H}_{2} \cdot{ }^{20}$ Noell and Hay, ${ }^{24 \mathrm{a}}$ however, reported a large HH stretch in their ab initio calculation, which suggests a large isotope effect. We have calculated the kinetic deuterium isotope effect $r$ at 300 K using the Bigeleisen equation ${ }^{51}$

$$
\begin{gather*}
r=\frac{k_{1}}{k_{2}} \frac{s_{2} s_{1}{ }^{*}}{s_{1} s_{2}{ }^{*}} \frac{\kappa_{2}}{\kappa_{1}}=R_{\mathrm{L}}^{*}\left(\prod_{i}^{3 N-7} R_{i}^{*}\right)\left(\prod_{i}^{3 N-6} R_{i}\right)^{-1}  \tag{1}\\
R_{\mathrm{L}}{ }^{*} \equiv \frac{\nu_{1 \mathrm{~L}}{ }^{*}}{\nu_{2 \mathrm{~L}}{ }^{*}}  \tag{2}\\
R_{i}^{\mathrm{a}} \equiv \frac{u_{1 i}{ }^{\mathrm{a}}}{u_{2 i}{ }^{\mathrm{a}}} \frac{\exp \left(-u_{1 i}^{\mathrm{a}} / 2\right)\left\{1-\exp \left(-u_{1 i}^{\mathrm{a}}\right)\right\}^{-1}}{\exp \left(-\mathrm{u}_{2 i}^{\mathrm{a}} / 2\right)\left\{1-\exp \left(-\mathrm{u}_{2 i}^{\mathrm{a}}\right)\right\}^{-1}}  \tag{3}\\
\mathrm{a}=\text { none or } \ddagger \\
u_{i} \equiv \hbar \nu_{i} / k T
\end{gather*}
$$

where $s$ is the symmetry number, $\kappa$ is the transmission coefficient, $\nu_{i}$ is the $i$ th vibrational frequency of the reactants, $\nu_{i}^{*}$ is that of the transition state, $\nu_{\mathrm{L}}{ }^{*}$ is the reaction-coordinate frequency, and subscripts 1 and 2 refer to the hydride and its deuterated systems, respectively.

The calculated isotope ratio $r$ is 1.48 , consistent with the experimental values of 1.2-1.3.9 Analysis of the isotope ratio is given in Table X where only factors giving $R$ values substantially different from unity are shown. The largest contributions to the isotope ratio $r$ are 1.43 from the weakening of the XX bond (X $=\mathrm{H}$ and D ) at the transition state and 1.31 from the reaction coordinate, i.e., the translational displacement of $X_{2}$ approaching Pt. Other modes contribute to moderate the effect of these two factors. In our calculation the contribution of the out-of-plane $\mathrm{B}_{2}$ bending has been assumed to be unity because of numerical errors introduced in calculating this small vibrational frequency. However, its real effect is expected to reduce the isotope ratio,
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(57) Hoyano, J. K.; McMaster, A. D.; Graham, W. A. G. J. Am. Chem. Soc. 1983, $105,7190$.
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Figure 4. Fully optimized geometry (in $\AA$ and deg) of the "transition state" leading directly to a trans compound found by forcing the $\mathrm{H}_{2}$ bond axis perpendicular to the $\operatorname{PPtP}$ plane with an overall $C_{2 v}$ symmetry. The state is unstable to the $\mathrm{H}_{2}$ rotation around the $C_{2 v}$ symmetry axis and therefore is not a true transition state. The smaller basis set.



Figure 5. Steric repulsion energy ( $\Delta E$ in $\mathrm{kcal} / \mathrm{mol}$ ) between two $\mathrm{PH}_{2}$ ( $i-\mathrm{Pr}$ ) molecules and the nearest HH distance ( $R_{\mathrm{HH}}$ in $\AA$ ) between two isopropyl groups. The smaller basis set.
making the calculated value even closer to the experimental value.
The transition-state geometry seems reasonable judged from the calculated low activation energy and small deuterium kinetic isotope effect. However, it may be suspected that Pt ECP parameters ${ }^{29}$ used in the present work may suffer a same defect as a similar ECP showed in an unrealistic attractive potential curve for the dichlorine molecule. ${ }^{59}$ In order to test ECP and basis set dependency, we fully optimized the geometry of the transition state using more reliable Pt ECP parameters and the basis set of Noell and $\mathrm{Hay}^{246}$ and have shown the result in the last column of Figure 2b. The transition state is apparently "an early one"; namely, the HH distance of $0.77 \AA$ is only a $5 \%$ increase from that of 0.731 $\AA$ in a free $\mathrm{H}_{2}$ molecule, and the HPt distance of $2.01 \AA$ is much longer than that of $1.51 \AA$ in the product cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$. The calculated activation energy is small, $10.0 \mathrm{kcal} / \mathrm{mol}$ at the RHF level, and is comparable with the present value of $5.2 \mathrm{kcal} / \mathrm{mol}$. From these results it can be concluded that the earliness and the low activation energy is not likely to be an artifact of the ECP parameters and basis functions. The rather "late" transition state obtained in a GVB calculation by Noell and Hay ${ }^{24 a}$ may be due to incomplete optimization of the geometry.
The transition state found at present work is apparently on the way to the cis product. We also searched a transition state leading directly to a trans product. In this search the $C_{2 v}$ symmetry is enforced by keeping the HH bond axis perpendicular to the PPtP plane. The "transition state" thus found, shown in Figure 4, has an activation energy of about $50 \mathrm{kcal} / \mathrm{mol}$ at the RHF level, much higher than the activation energy of several kilocalories per mole
(59) Hay, P. J.; Wadt, W. R.; Kahn, L. R. J. Chem. Phys. 1978, 68, 3059.

## Scheme II


of the path leading to a cis product. This "transition state", however, is found to have two negative force constants, one for the reaction and another for the HH rotation around the $C_{2 v}$ axis, and, therefore, is not a true transition state. The reactants trying to follow this path will be deflected down toward the true transition state that leads to the cis product. We conclude that the transition state that directly leads to the trans product does not exist in the $\mathrm{H}_{2}$ oxidative addition. This is consistent with the correlation diagram discussed in the preceding section that the trans addition of $\mathrm{H}_{2}$ is symmetry forbidden.

There occurs a question why only the trans product is detected by experiment. Before discussing this problem, we must consider the change of relative stability of cis and trans products caused by a steric repulsion between phosphine ligands. We note that all the compounds used in the experiments have bulky phosphines such as $\mathrm{PPh}_{3}$ and $\mathrm{P}(t-\mathrm{Bu})_{3}$ rather than $\mathrm{PH}_{3}$ in our model system. We will estimate the steric repulsion effect by calculating the RHF interaction energy between two isopropylphosphines, $\mathrm{PH}_{2}(i-\mathrm{Pr})$. As shown in Figure 5, two phosphorus atoms of isopropylphosphines are placed at a distance of $2.194 \AA$, the PtP bond length in $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, from the empty origin Z and the PZP angle is changed between $120^{\circ}$ and $180^{\circ}$. In the calculation the structure of $\mathrm{PH}_{2}(i-\mathrm{Pr})$ is fixed at an optimized geometry of a free state within $C_{s}$ symmetry. The calculated interaction energy $\Delta E$ and the shortest isopropyl H-H distance $R_{\mathrm{HH}}$ are shown in Figure 5. One finds that at the PZP angle of about $140^{\circ} R_{\mathrm{HH}}$ is reduced to the sum of the hydrogen van der Waals radii, $2.4 \AA$, and the steric repulsion becomes very serious. In real Pt diphosphine compounds bulky phosphines will readjust their geometry to reduce the steric repulsion between them. Therefore the interligand angle where the steric repulsion becomes severe should be smaller than that obtained from the present model calculation. From these results we deduce that the steric repulsion between bulky phcsphines is not serious at the transition state, where the PPtP angle is $148^{\circ}$. However, in cis- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PR}_{3}\right)_{2}$ the repulsion is expected to be very large, making this complex much less stable than the corresponding trans complex. We note that Noell and Hay estimated the steric repulsion to be $7 \mathrm{kcal} / \mathrm{mol}$ at the PPtP interligand angle of $100^{\circ}$, by using trimethylphosphine ligands.

On the basis of these discussions, one may imagine potential energy profiles shown in Scheme II involving only one early and cis-leading transition state. In the scheme the $x$ direction stands for an in-plane $\mathrm{H}_{2}$ approach and the $y$ direction a perpendicular approach to the PPtP plane, dots specify the reactant and products, T means the transition state, and X represents the unaccessible "transition state". In Scheme IIa where the interligand steric repulsion is negligible, a deep valley path efficiently leads the system to the cis product. In Scheme IIb where bulky phosphines are used, the system still goes through the same early transition state where the steric repulsion is not yet serious, and, as it approaches the product, the steric repulsion has a good chance to push the system away from the cis product and lead to the less hindered trans product. There are three possible mechanisms of trans product formation: (1) the four-coordinate mechanism as shown in Scheme IIb, (2) the three-coordinate mechanism in which one of the bulky phosphines is lost, so that the steric repulsion between them is released, and (3) the five-coordinate mechanism. Experimentally the oxidative addition of $\mathrm{H}_{2}$ takes place in benzene, toluene, and $n$-hexane, ${ }^{4,5}$ less coordinating solvents; and therefore the first two mechanisms are more plausible than the last. Although the present calculation does not show which of the first

Table XI. Analysis of Kinetic Deuterium Isotope Effect at 250 K for the Reductive Elimination cis- $\mathrm{Pt}(\mathrm{X})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2} \rightarrow \mathrm{CH}_{3} \mathrm{X}+$ $\underline{\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}(\mathrm{X}=\mathrm{H} \text { and } \mathrm{D})^{a}}$

| mode (sym) | reactant |  | transitn state |  | $\begin{gathered} \prod_{i} \\ R_{i}^{*} R_{i}^{-1} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\nu,{ }^{\text {b }} \mathrm{cm}^{-1}$ | $R_{i}$ | $\nu,{ }^{\text {b }} \mathrm{cm}^{-1}$ | $R_{i}{ }^{\text { }}$ |  |
| XPtC in-plane bending ( $\mathrm{A}^{\prime}$ ) | 1150 | 0.55 | 1030i | 1.27 |  |
|  | 680 | 0.98 |  |  | 1.80 |
|  |  |  | 1090 | 0.82 |  |
|  |  |  | 1430 | 0.94 |  |
| XPt stretching ( $\mathrm{A}^{\prime}$ ) | 2160 | 0.23 | 2170 | 0.34 | 1.16 |
|  |  |  | 1630 | 0.79 |  |
| XPtC out-of-plane bending ( $\mathrm{A}^{\prime \prime}$ ) | 780 | 0.85 | 1080 | 0.80 |  |
|  | 490 | 0.91 | 430 | 0.93 | 0.94 |
|  |  |  | 580 | 0.99 |  |
|  |  |  | 620 | 0.99 |  |
| total ${ }^{\text {c }}$ |  | 0.096 |  | 0.19 | 1.97 |

${ }^{a}$ Only factors contributing significantly to the isotope effect are listed. ${ }^{b}$ The frequencies are those in hydride compounds $\mathrm{Pt}(\mathrm{H})(\mathrm{C}$ $\left.\mathrm{H}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$. ${ }^{\text {c }}$ Total product of contributions of all the modes where for $\mathrm{PH}_{3}$ rotations around PtP bonds, which have near-zero frequencies, their $R_{\mathrm{i}}$ and $R_{i}{ }^{*}$ are assumed to be 1 .
two actually takes place, it can be said that the steric repulsion between the bulky phosphines is one of the important keys of the reaction, and experimental studies with less bulky phosphines would shed more light in elucidating the detailed mechanism.
(B) Transition State and Kinetic Isotope Effect in the Reductive Elimination of $\mathrm{CH}_{4}$ from cis $-\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$. The reduc-tive-elimination reaction of $\mathrm{CH}_{4}$ from $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PPh}_{3}\right)_{2}$ has been found by Abis, Sen, and Halpern. ${ }^{8}$ The reaction rate was unaffected by the presence of excess triphenylphosphine, suggesting strongly that the elimination proceeds through a four-coordinate transition state. A large deuterium kinetic isotope ratio of 3.3 was observed for the decomposition reaction of $\mathrm{Pt}(\mathrm{D})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$.
We have searched the transition state for the reaction $\mathrm{CH}_{4}+$ $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2} \rightarrow \mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, the oxidative addition of $\mathrm{CH}_{4}$ to $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$, and the reductive elimination of $\mathrm{CH}_{4}$ from cis- Pt $(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, the arrows shown in Figure 3 indicate the reaction coordinate vectors. Normal-mode frequencies are given in Table IV, where the imaginary frequency of $1030 \mathrm{i} \mathrm{cm}^{-1}$ corresponds to the reaction coordinate.
The conformation of the transition state has the following features: (1) the $\operatorname{PPtP}$ angle of $119^{\circ}$ is much smaller than that of $148^{\circ}$ at the transition state for the $\mathrm{H}_{2}$ oxidative addition, (2) the PtC distance of $2.27 \AA$ is $14 \%$ longer than that of $2.00 \AA$ in the reactant cis- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, (3) the CH distance of 1.35 $\AA$ is $25 \%$ longer than that of $1.08 \AA$ in the product $\mathrm{CH}_{4}$, and (4) the PtH distance of $1.59 \AA$ is only a $4 \%$ increase from that of 1.52 $\AA$ in cis- $\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$. The first three features indicate that the transition state is located at a midpoint on the potential energy hypersurface between cis $-\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}+\mathrm{CH}_{4}$. This is in contrast with the case of $\mathrm{H}_{2}$, where the transition state is located near $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}+\mathrm{H}_{2}$. The last feature shows that, though the reaction is concerted, two bonds do not break simultaneously in the elimination reaction. The weaker $\mathrm{C}-\mathrm{Pt}$ bond is substantially broken before the $\mathrm{H}-\mathrm{Pt}$ bond is replaced by a new $\mathrm{C}-\mathrm{H}$ bond. In the reverse oxidative-addition reaction, of course, the $\mathrm{H}-\mathrm{Pt}$ bond is formed in an early stage of reaction.

Interestingly the reaction coordinate at the transition state in Figure 3 has characteristics of HPtC bending rather than PtH stretching. It is important, therefore, to find out whether this transition state is consistent with the large observed kinetic deuterium isotope ratio of $3.3 .^{8}$ We have calculated and analyzed the isotope effect by using the Bigeleisen equation ${ }^{51}$ and have given the results in Table XI. The calculated ratio of 2.0 is in reasonable agreement with the experimental value. The largest contribution to the calculated ratio is mainly due to 1.80 of the HPtC bending mode, as expected from the transition-state reaction coordinate.
We have also calculated the deuterium kinetic isotope ratio in the reverse reaction, the oxidative addition, of $\mathrm{CH}_{3} \mathrm{D}$ and $\mathrm{CD}_{4}$ to $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ and obtained values of $6.7(250 \mathrm{~K})$ and $5.0(300 \mathrm{~K})$ for $\mathrm{CH}_{3} \mathrm{D}$ and $17.0(250 \mathrm{~K})$ and $12.0(300 \mathrm{~K})$ for $\mathrm{CD}_{4}$. The $\mathrm{CH}_{4}$

Oxidative Addition and Reductive Elimination Reactions

Table XII. Energy Decomposition Analysis ${ }^{a}$ of the Activation Energy of the Oxidative Addition of $\mathrm{H}_{2}$ to $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ (in $\mathrm{kcal} / \mathrm{mol}$ )

| deformation |  |
| :--- | ---: |
| $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ | 3.1 |
| $\mathrm{H}_{2}$ | 0.3 |
| sum | 3.4 |
| interaction |  |
| ES + EX | 16.9 |
| BCTPLX | -7.2 |
| FCTPLX | -10.7 |
| residue | 2.8 |
| sum | 1.8 |
| total | 5.2 |

${ }^{a}$ With the larger basis set.
addition reaction has not been observed; however, the calculated isotope ratio will be useful to the future kinetic study of CH bond addition reactions such as orthometalations. ${ }^{52}$

In order to discuss the reaction mechanism in detail, the effect of the steric repulsion between bulky phosphines has to be considered as they are usually used in experiments. As discussed in the previous section, the repulsion will be serious in the cis compound, where the PPtP angle is about $120^{\circ}$. This destabilization would make the oxidative addition endothermic, which explains why the reaction has not been observed, and the reverse reductive elimination becomes exothermic. Accompanied with this change, the position of the transition state on the reaction coordinate should move further toward the cis compound, which is in contrast with the less affected transition state in the $\mathrm{H}_{2}$ reaction. This implies that the possibility of unimolecular isomerization of cis hydride methyl compound $\operatorname{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PR}_{3}\right)_{2}$ is small, because the proximity of this transition state would lead to decomposition. The change of location of the transition state will also modify the reaction coordinate vector. However, the HPtP bending mode will remain to be an important component of the reaction coordinate, since the PtH distance at the transition state is almost the same as that in the cis compound. Therefore, a large isotope due to the HPtP bending mode is expected to remain.
(C) Important Factors for Reaction. Two possible ways of a bond activation by metal compounds are known: a forward electron transfer from a bonding orbital of a breaking bond into vacant metal orbitals ${ }^{53}$ and a backward electron transfer from occupied metal d orbitals into an antibonding orbital. ${ }^{54,55}$ Dedieu et al. ${ }^{20}$ and Lauher et al. ${ }^{22}$ have pointed out from theoretical studies that the backward electron transfer is important in the stabilization of the transition state of the $\mathrm{H}_{2}$ addition reaction. In order to examine the charge redistribution, we have carried out Mulliken population analysis at the transition states of $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}+\mathrm{HX}(\mathrm{X}$ $=\mathrm{H}$ and $\mathrm{CH}_{3}$ ), as given in columns 8 and 9 of Table I. One finds at the $\mathrm{H}_{2}$ addition transition state that the population 1.82 of Pt $\mathrm{d}_{y z}$ is smaller than 1.88 of the reactant and that the population 0.17 of $\mathrm{Pt} \mathrm{p}_{y}$ is larger than 0.02 of the reactant. Though columns 2 and 5 of Table I show that just a bend of $\mathrm{PH}_{3}$ ligands causes similar changes in the population, their magnitude is much smaller than the change found at the $\mathrm{H}_{2}$ addition transition state. The above changes suggest both the backward electron transfer from the $\mathrm{Pt} \mathrm{d}_{y z}$ orbital to $\mathrm{H}_{2} \sigma^{*}$ and the forward transfer from $\mathrm{H}_{2} \sigma$ to $\mathrm{Pt} \mathrm{p}_{y}$ are important at the transition state. This will soon be confirmed by the energy analysis.

In order to investigate further the nature of electronic interaction, we have carried out the energy decomposition analysis (EDA) ${ }^{56}$ at the transition state of the $\mathrm{H}_{2}$ addition reaction. In the EDA, the RHF interaction energy is decomposed into the deformation energy of fragments from their equilibrium geometries and the electronic interaction energy among fragments. The latter energy is further decomposed into the sum of electrostatic and exchange interaction ( $\mathrm{ES}+\mathrm{EX}$ ) between fragments, the forward and the backward charge-transfer interaction, FCTPLX and BCTPLX, respectively, and the residual interactions. In the present analysis $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ is divided into $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{H}_{2}$ fragments and the interaction between them is investigated. Results are given in Table XII, where a negative number means


Figure 6. PtP bond strength (BS) and orbital energy $\epsilon$ of higher occupied orbitals in $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ as functions of the PtP angle. The BS is the energy difference between $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)+\mathrm{PH}_{3}$ at each PPtP angle. The MO symmetry in the lower figure is the same as that in Figure 1. The smaller basis set.


Figure 7. PtH bond strength ( BS ), PtH and PtP bond distances, and HPtH bond angles, $\psi$ in cis- $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ as a function of the PPtP angle, $\theta$. The BS is a half of the energy difference between $\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}+\mathrm{H}+\mathrm{H}$. The bond distances are relative to those at the angle $\theta$ of $100^{\circ}$. The smaller basis set.
a stabilization.
In the $\mathrm{H}_{2}$ addition reaction the energy required to deform two fragments from the equilibrium geometry to the geometry taken at the transition state is only $3.4 \mathrm{kcal} / \mathrm{mol}$, typical of an "early transition state". The electronic interaction energy is $1.8 \mathrm{kcal} / \mathrm{mol}$ repulsive, where the electrostatic and exchange destabilization of $17 \mathrm{kcal} / \mathrm{mol}$ is nearly compensated by the forward charge-transfer (FCTPLX) stabilization of $11 \mathrm{kcal} / \mathrm{mol}$ as well as the backward charge-transfer (BCTPLX) stabilization of $7 \mathrm{kcal} / \mathrm{mol}$. The "oxidation" of the metal will occur at some later stage of the reaction when the HH bond is actually broken.

The large backward charge-transfer interaction, BCTPLX, is consistent with the decreased total gross population in metal d and $\mathrm{p}_{z}$ orbitals. The large forward electron-transfer, FCTPLX, stabilization is reflected in the increased $d_{y^{2}}$ and $p_{y}$ population at the transition state. This FCTPLX interaction indicates that the


Figure 8. The relation between the bond strength of PtH in cis- $\mathrm{Pt}-$ $(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and the $\mathrm{d}_{y 2}$ orbital energy in $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$. The numbers in the figure are the value of the PPtP angle. The smaller basis set.
metal is in a sense "reduced" as well as "oxidized" at an early stage of the reaction.

In order to control the reactivity of two-coordinate $\operatorname{Pt}(0)$ and $\operatorname{Pd}(0)$ diphosphine compounds, Otsuka et al. suggested to consider the interligand angle as well as the steric size and basicity of phosphine ligands. $4.5,17$ We have calculated the PtP bond strength and orbital levels in $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ and PtH bond strength in cis- Pt $(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ as functions of the PPtP interligand angle and show the results in Figures 6 and 7, respectively.

In the two-coordinate $\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ the PtP bond strength is positive in the range of the $\operatorname{PPtP}$ angles from about $80^{\circ}$ to $180^{\circ}$ (Figure 6), which suggests that Pt-chelating phosphine complexes can exist at these PPtP angles. At PPtP angles less than $80^{\circ}$ one of the phosphines will probably dissociate. The change of the PtP bond strength is estimated to be about $10 \mathrm{kcal} / \mathrm{mol}$ for a decrease in the $\operatorname{PPtP}$ angle of $10^{\circ}$ in the range between $100^{\circ}$ and $140^{\circ}$. Among the higher occupied MO's only a $b_{1}$ MO consisting mainly of the $\mathrm{Pt} \mathrm{d}_{y z}$ orbital is destabilized as the PPtP angle decreases. Electrons in this MO transfer to the reactant such as $\mathrm{H}_{2}$ in the oxidative addition. Therefore, a chelating phosphine having a small PPtP angle effectively raises the reactivity of Pt electrons.

The dependence of the PtH bond strength and three geometrical parameters, PtH and PtP distances and the HPtH angle, on the PPtP interligand angle in cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ is shown in Figure 7. For a decrease of the PPtP angle by $10^{\circ}$, the PtH bond strength increases by about $10 \mathrm{kcal} / \mathrm{mol}$, a dependence opposite to that of the $\operatorname{PtP}$ bond strength. We note that in the possible range of $\theta$, the change in the PtH and PtP bond distance is small. Relationships among the PtH bond strength, the $\mathrm{d}_{y z}$ orbital energy, and the PPtP angle are depicted in Figure 8. The change of the PtH bond strength in cis $-\mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ correlates very well with the destabilization of the $\mathrm{d}_{y z}$ orbital.
The oxidative addition of methane to transition-metal compounds has been a challenging problem for experimentalists. Recently Graham and Bergman have independently succeeded in carrying out the reaction by using iridium compounds. ${ }^{57,58}$ The presumed reaction intermediates $\mathrm{CpIrX}\left(\mathrm{X}=\mathrm{CO}, \mathrm{PMe}_{3}\right)$ are structurally unsaturated and have open-shell electrons (a) or a vacant 5d orbital (b). They are probably highly reactive, as a


vacant $d$ orbital has been found to strongly activate a CH bond
and makes the activation energy small for bond cleavage. ${ }^{60}$ Although Pt diphosphine does not have a vacant $d$ orbital, according to the above semiquantitative analysis there may be a possibility of $\mathrm{CH}_{4}$ oxidative addition to Pt with a chelating phosphine such as $\mathrm{Pt}\left[\mathrm{R}_{2} \mathrm{P}\left(\mathrm{CH}_{2}\right)_{3} \mathrm{PR}_{2}\right](\mathrm{R}=t-\mathrm{Bu})$ where the PPtP angle has been observed to be $102.6^{\circ}, 5$ although the dimerization such as observed by Yoshida et al. ${ }^{5}$ has to be prevented in the first place.

## Summary

Potential energy hypersurfaces of the reactions $\mathrm{H}_{2}+\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}$ $\rightarrow \mathrm{Pt}(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}$ and $\mathrm{CH}_{4}+\mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2} \rightarrow \mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$ are investigated theoretically by fully optimizing the transition state as well as the equilibrium geometries. These two hypersurfaces show the following contrasts. The $\mathrm{H}_{2}$ addition reaction is exothermic, while the corresponding $\mathrm{CH}_{4}$ addition reaction is almost thermoneutral. The transition state of the former reaction is in an early stage with the PPtP angle of $148^{\circ}$, whereas that of the latter is at the midpoint with the angle of $119^{\circ}$. The activation barrier is only $8 \mathrm{kcal} / \mathrm{mol}$ in the former, while it is $28 \mathrm{kcal} / \mathrm{mol}$ in the latter. The reaction coordinate vectors consist mainly of the $\mathrm{H}_{2}$ translational mode in the former and of the HPtP bending mode in the latter. The large hydrogen isotope effect observed for the $\mathrm{CH}_{4}$ reductive elimination is consistent with the calculated transition state.

These differences in the potential energy hypersurfaces can be explained by the differences between the PtH and PtC bond strengths that are estimated to be $61-70$ and $34-41 \mathrm{kcal} / \mathrm{mol}$, respectively. The larger PtH bond strength results in a larger exothermicity, an earlier transition state, and hence a lower activation barrier.

The effect of the steric repulsion between bulky phosphines is expected to be different between $\mathrm{H}_{2}$ and $\mathrm{CH}_{4}$. In the $\mathrm{H}_{2}$ addition reaction the effect would be negligible at the cis transition state but destabilizes the cis product. Therefore the system would easily go through the cis transition state and then deviate from the sterically undesirable cis path to reach the less destabilized trans product. In the $\mathrm{CH}_{4}$ addition reaction the effect would be large even at the transition state, making the reaction all the more difficult.

The decrease of PPtP interligand angle activates efficiently and selectively one of the metal $d$ orbitals. The activated energy is calculated to be about $10 \mathrm{kcal} / \mathrm{mol}$ per $10^{\circ}$ change in the PPtP angle in the range from $100^{\circ}$ to $140^{\circ}$.

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Registry No. cis $-\mathrm{Pt}(\mathrm{H})\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}$, 79232-17-0; trans- $\mathrm{Pt}(\mathrm{H})$ $\left(\mathrm{CH}_{3}\right)\left(\mathrm{PH}_{3}\right)_{2}, 92541-64-5 ;$ cis $-\mathrm{Pt}\left(\mathrm{H}_{2}\left(\mathrm{PH}_{3}\right)_{2}, 76832-29-6 ;\right.$ trans $-\mathrm{Pt}-$ $(\mathrm{H})_{2}\left(\mathrm{PH}_{3}\right)_{2}, 76830-84-7 ; \mathrm{Pt}\left(\mathrm{PH}_{3}\right)_{2}, 76830-85-8 ; \mathrm{CH}_{4}, 74-82-8 ; \mathrm{H}_{2}$, 1333-74-0; $D_{2}$, 7782-39-0.
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